

Diffusively coupled networks

Paul Van den Hof

Doctoral School Lyon, France, 11-12 April 2024

www.sysdynet.eu www.pvandenhof.nl p.m.j.vandenhof@tue.nl



p.m.j.vandenhof@



Diffusively coupled networks

Back to the basics of physical interconnections

In connecting physical systems, there is often no predetermined direction of information^[1]



Example: resistor / spring connection in electrical / mechanical system:



Difference of node signals drives the interaction: diffusive coupling

Diffusively coupled physical network



Equation for node *j*:

$$M_j \ddot{w}_j(t) + D_{j0} \dot{w}_j(t) + \sum_{k \neq j} D_{jk} (\dot{w}_j(t) - \dot{w}_k(t)) + K_{j0} w_j(t) + \sum_{k \neq j} K_{jk} (w_j(t) - w_k(t)) = u_j(t),$$



Mass-spring-damper system

- Masses M_j
- Springs K_{jk}
- Dampers D_{jk}
- Input u_j



$$\begin{bmatrix} M_{1} & M_{2} & \\ & M_{3} \end{bmatrix} \begin{bmatrix} \ddot{w}_{1} \\ \ddot{w}_{2} \\ \ddot{w}_{3} \end{bmatrix} + \begin{bmatrix} 0 & D_{20} & \\ & D_{20} & 0 \end{bmatrix} \begin{bmatrix} \ddot{w}_{1} \\ \ddot{w}_{2} \\ \dot{w}_{3} \end{bmatrix} + \begin{bmatrix} K_{10} & 0 & \\ & 0 & 0 \end{bmatrix} \begin{bmatrix} w_{1} \\ w_{2} \\ w_{3} \end{bmatrix} \\ + \begin{bmatrix} D_{13} & 0 & -D_{13} \\ 0 & D_{23} & -D_{23} \\ -D_{13} & -D_{23} & D_{13} + D_{23} \end{bmatrix} \begin{bmatrix} \dot{w}_{1} \\ \dot{w}_{2} \\ \dot{w}_{3} \end{bmatrix} + \begin{bmatrix} K_{12} + K_{13} & -K_{12} & -K_{13} \\ -K_{12} & K_{12} & 0 \\ -K_{13} & 0 & K_{13} \end{bmatrix} \begin{bmatrix} w_{1} \\ w_{2} \\ w_{3} \end{bmatrix} = \begin{bmatrix} 0 \\ u_{2} \\ 0 \end{bmatrix} \\ \begin{bmatrix} X(p) \\ diagonal \end{bmatrix} + \begin{bmatrix} Y(p) \\ Laplacian \end{bmatrix} w(t) = u(t) \qquad X(p), Y(p) \text{ polynomial } p = \frac{d}{dt} \end{bmatrix}$$

Mass-spring-damper system



This fully fits in the earlier module representation:

$$w(t) = Gw(t) + \underbrace{Rr(t) + He(t)}_{Q^{-1}(p)u(t)}$$

with the additional condition that:

 $G(p) = Q(p)^{-1}P(p)$ Q(p), P(p) polynomial P(p) symmetric, Q(p) diagonal



Module representation

Consequences for node interactions:



- Node interactions come in pairs of modules
- Where numerators are the same

Framework for network identification remains the same

• Symmetry can be incorporated in identifiability/identification

Polynomial representation

More attractive: stay within the polynomial domain (discrete-time now)

$$[\underbrace{Q(q^{-1})}_{}-\underbrace{P(q^{-1})}_{}] w(t) = u(t)$$

diagonal

hollow&symmetric

$$A(q^{-1})w(t) = \underbrace{B(q^{-1})r(t) + v(t)}_{u(t)}$$

with $A(q^{-1})$ symmetric and nonmonic i.e. $A(q^{-1}) = A_0 + A_1 q^{-1} + \cdots + A_n q^{-n}$ with $A_0 \neq I$

Network identifiability^[1]

New analysis, based on $T_{wr}(q)$ only (noise discarded because of algebraic loops):

 $A(q^{-1})w(t) = B(q^{-1})r(t)$ $\Pi(q^{-1}) \left[A(q^{-1})w(t) = B(q^{-1})r(t) \right]$

Identifiability conditions:

- At least 1 excitation signal r(t) present
- $A(q^{-1})$ and $B(q^{-1})$ left coprime
- diagonality constraint on $[A_0 \cdots A_n \ B_0 \cdots B_n]$
- $A(q^{-1})$ symmetric
- 1 parametric constraint in $A(q^{-1})$ or $B(q^{-1})$ 1

- $B(q^{-1})$ present
- $\Pi(q^{-1})$ unimodular
- **Π** diagonal
- $\Pi = \alpha I$
- $\Pi = I$

Polynomial representation - identifiability

- Identifiability conditions are strongly relaxed (compared to module framework) in terms of number of excitation signals required.
- Diffusive couplings strongly limit the **degrees of freedom** in the network model
- Identification algorithms are available for both full network^[1] and local identification^[2].

10



Summary diffusively coupled networks

- Interesting class of models, not extensively studied in identification
- Non-directed graphs
- Adhering to physical interconnections
- Framework is fit for representing combined networks (combining physical bi-directional links, and cyber uni-directional links)^{[1].}



The end